

# The Body Doping Concentration Effects on the CNTFETs Characteristics

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**Abstract**— This paper presents the effects of body doping concentration on a carbon nanotube field-effect transistor (CNTFET) with symmetrical coaxial geometry structure, that transistor has a wrapped gate dielectric around the nanotube. At first, Gordon Moore's law is expressed briefly. Then, the limitations and problems after 10nm channel lengths of the traditional metal-oxide-semiconductor field-effect transistors (MOSFETs) scaling are introduced and the structure of carbon nanotube field-effect transistor as a solution for limitations is illustrated. The simulation process is performed, then the effect of body doping concentration is investigated on the current-voltage characteristics, energy profile and electron density of transistor. The simulation is based on a self-consistent solution of the two dimensions Poisson equation and Schrodinger solver. Results show that with increasing body doping concentration level, the transistor output current, energy profile and electron density of transistor will increase and on/off current ratio will decrease.

**Index** Carbon Nanotube Field Effect Transistor (CNTFET); body doping; self-consistent solution; Schrödinger equation.

## I. INTRODUCTION

IN 1965 Gordon Moore predicted that the number of transistors on a chip will double every 18 months. In other words, the dimensions of devices will decrease to half [1]. Scaling down of devices is necessary for advanced electronic technology. Traditional MOSFETs have many important limitations and problems after 10nm channel lengths [2]. According to ITRS, MOSFETs will reach their limits in 2020 [3]. The limitations are electron tunneling through short channels and thin insulator films, the associated leakage currents, passive power dissipations, short channel effects, increasing cost of the fabrication process, variations in device the nanotube. The reason for this kind of structure selection is the cylindrical nature of CNTs. This structure provides the best channel controlling by gate field, as sequence the minimum

structure and doping process. Therefore new inventions are needed [4]. The semiconductor industry is looking for new materials for replacing silicon in electronic devices [5]. One of the most promising building blocks for the future of nanoelectronic devices is Carbon Nanotubes (CNT).

Iijima discovered Carbon nanotubes (CNTs) as a by-product in 1991 in NEC Corporation [6]. CNTs are cylindrical of carbon sheets that have many numbers of unique features and special properties such as high electrical conductivity, heat conductivity, metallic or semiconducting behavior, ballistic electron transport, huge current carrying capacity, high mobility, compatibility with high dielectric constants (K) and small diameters that are suitable for the basic building blocks in future nanoelectronic applications. In order to use this excellent property of CNTs, scientists have fabricated Carbon Nanotube Field Effect Transistors (CNTFETs) based on single-walled semiconducting CNTs as channel material, that fabricated for the first time in 1998, hence CNTFETs are one of the novel nanoelectronic devices that overcome MOSFETs limitations. In order to study the performance of CNTFETs, the self-consistent solution of the Poisson and Schrodinger equations is used, then the body doping concentration effect on the CNTFETs characteristics has been investigated [7].

## II. TWO-DIMENSIONAL SIMULATION OF CARBON NANOTUBE FIELD-EFFECT TRANSISTORS

In this section, electron transport simulation in a semiconducting nanotube that using in the CNTFET structure is described. For the accurate simulation of CNTFETs self-consistent solution of the Poisson and Schrodinger equations is used. The CNTFET with Coaxial geometry structure is used in this research, the transistor has a wrapped gate dielectric around channel length. Fig.1 shows the structure of CNTFET that is used in this work.

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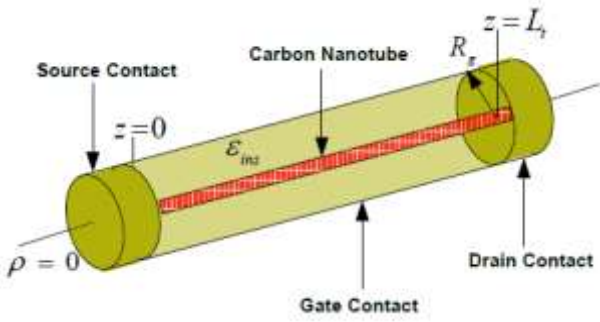


Fig.1: Coaxial CNTFET model geometry.

Drain and source contacts have the same radius and length with gate contact. The nanotube in this device structure is surrounded by an insulator with relative permittivity  $\epsilon_{ins}$ . The source and drain contacts terminate the ends of the device.  $R_g$  is the gate radius,  $R_t$  is nanotube radius, the insulator thickness  $t_{ins}=R_g-R_t$  and the nanotube length is  $L_t$ . In this system due to its symmetrical geometry, Poisson's equation, restricted to just two dimensions [8]–[10].

$$\frac{\partial^2 V}{\partial \rho^2} + \frac{1}{\rho} \frac{\partial V}{\partial \rho} + \frac{\partial^2 V}{\partial z^2} = -\frac{Q}{\epsilon} \quad (1)$$

Where the  $Q$  is charge density and  $V(\rho, z)$  is the potential within the outer cylinder. The charge distribution on the surface of the nanotube is obtained by solving the Poisson equation coupled with the time-independent Schrödinger equation:

$$\frac{\partial^2 \Psi}{\partial z^2} = -\frac{\gamma m^*}{\hbar^2} (E - U) \Psi = -K^2 \Psi \quad (2)$$

Where  $\Psi(z, E)$  is the wave function of a carrier with total energy  $E$  and effective mass  $m^*$ , traveling in a region with local effective potential  $U(z)$ . (1) solving is required with  $Q=Q(V)$ . Convergence for this non-linear system is achieved when:

$$V_{k+1} = V_k - \alpha L^{-1} r_k \quad (3)$$

$$r_k = LV_k + Q(V_k) \quad (4)$$

Where  $r_k$  is the residual of the  $k$ -th iteration,  $\alpha$  is a damping parameter, and  $L$  represents the differential operator allowing (1) to be written as:

$$LV = -Q \quad (5)$$

The energy of the nanotube vacuum level is calculated by:

$$E_{vac}(z) = -qV_{GS}(z) \quad (6)$$

The potential of electrons and holes in the nanotube are given by (7) and (8).

$$U_e(z) = E_{vac}(z) - \chi_{CN} \quad (7)$$

$$U_h(z) = -U_e(z) + E_g \quad (8)$$

Where  $E_g$  is bandgap of nanotube and  $\chi_{CN}$  is electron affinity. The charge density is calculated by:

$$Q = \frac{q(p - n) \delta(\rho - R_t)}{\rho} \quad (9)$$

Where  $\delta(0)/\rho$  is Dirac delta function in cylindrical coordinates,  $p$  and  $n$  are electron and hole density that is achieved from the Schrodinger equation. The solution of Schrodinger equations is performed with using the scattering-matrix method. The total of the system is divided into  $N$  equal sections. For each section, the electron wave function is introduced by:

$$\Psi_i = A_i e^{ik_i z} + B_i e^{-ik_i z} \quad i = 1, 2, \dots, N \quad (10)$$

Where  $k$  is wave vector,  $A$ ,  $B$  are wave function amplitude. Boundary conditions between  $i$ ,  $i+1$  sections of the system are:

$$\Psi_i = \Psi_{i+1} \quad (11)$$

$$\frac{\partial \Psi_i}{\partial z} = \frac{\partial \Psi_{i+1}}{\partial z} \quad (12)$$

According to scattering-matrix theory, the relationship between

$$\begin{bmatrix} A_i \\ B_i \end{bmatrix}, \begin{bmatrix} A_{i+1} \\ B_{i+1} \end{bmatrix} \text{ is:} \quad \begin{bmatrix} A_{i+1} \\ B_{i+1} \end{bmatrix} = S_i \begin{bmatrix} A_i \\ B_i \end{bmatrix} \quad (13)$$

Where  $S_i$  is scattering-matrix,  $\Psi_1$  and  $\Psi_N$  are wave functions at the source and drain contacts. For simulation, the initial charge density inside the device usually is assumed zero. The electron potential inside the device is achieved by solving Poisson's equation. Then with using the scattering matrix method for solving the Schrodinger equation, the net charge density in the device is computed. Then these results used to solve the Poisson equation again and these steps are repeated until between two steps convergence achieved.

### III. SIMULATION RESULTS

In this section, the results of the simulation based on the CNTFET theory described in section 2, is considered. The carbon nanotube used in this CNTFET structure is a zig-zag nanotube with chirality (16,0). The diameter of the nanotube is 0.63nm, and the corresponding bandgap is 0.62eV. Nanotube length is 20nm, also source and drain contacts lengths are 20nm, the dielectric is hafnium with thickness,  $t_{ins}=2$ nm, and  $\epsilon_{ins}=25$ .

Fig.2 and 3 show the  $I_{DS}-V_D$  and  $I_{DS}-V_G$  characteristics of CNTFET with mentioned parameters in different body doping concentrations. The Body doping concentration is changed from  $N_b=0\text{cm}^{-3}$  to  $N_b=1.3 \times 10^9 \text{cm}^{-3}$  and results are shown in Fig. 2-7. When the concentration is below than  $N_b=1 \times 10^7 \text{cm}^{-3}$ , Output characteristics changes are not noticeable. According to simulation results, maximum body doping concentration that we can apply to this CNTFET is  $N_b=1 \times 10^9 \text{cm}^{-3}$  When the doping concentration is higher than this value, simulation results are not agreeable.

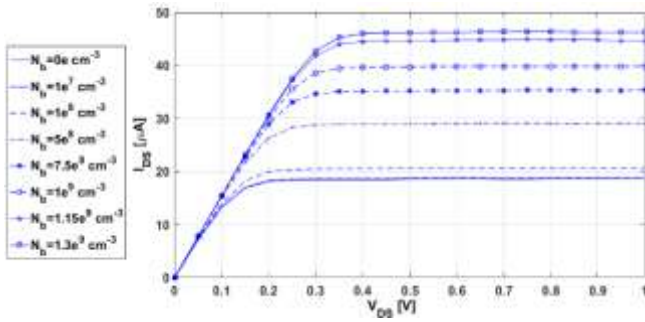


Fig. 2:  $I_{DS}$ - $V_D$  characteristics of CNTFET for different body doping concentrations, at  $V_{GS}=0.5V$

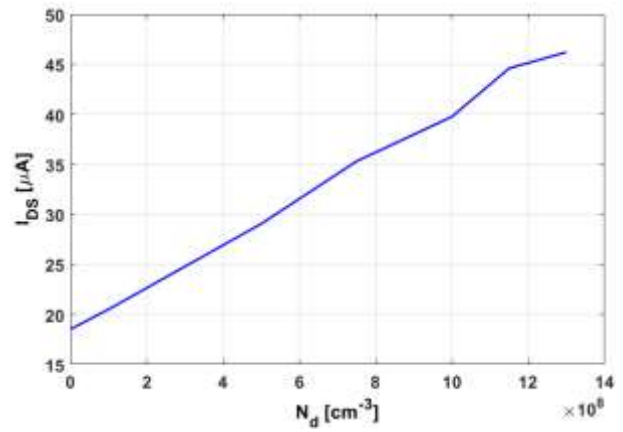


Fig.4: Output current versus body doping concentrations.

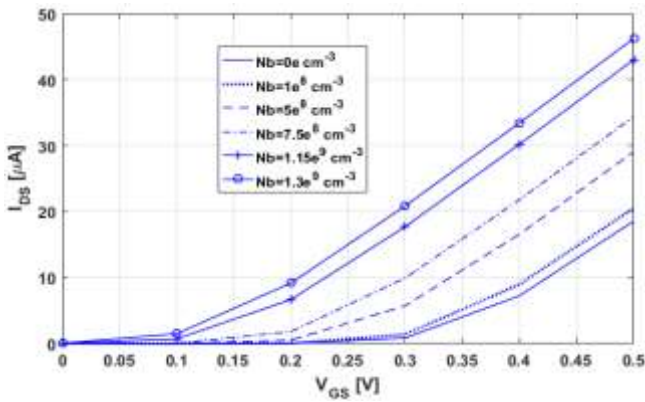


Fig. 3:  $I_{DS}$ - $V_G$  characteristics of CNTFET for different body doping concentrations, at  $V_{DS}=1V$

According to Fig.4, with increasing body doping concentration, the output current increases. In this CNTFET structure, there is a barrier region near the source/drain contacts and nanotube channel. These barriers determine the number of electrons entering the channel. The height of barriers is modulated by the body doping concentration. With increasing the body doping level, there is a consequent decrease in the barrier height. Hence, more carriers flow through the channel and output current will increase. In Fig.4, the Output current is shown as a function of body doping concentration.

According to Fig.5, with increasing body doping concentration, the energy profile of CNTFET increases. More doping concentration allows more propagation in the channel.

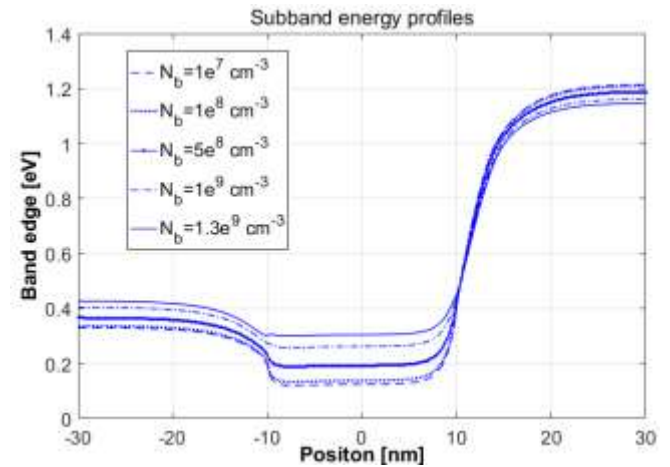


Fig. 5: Energy profile of CNTFET for different body doping concentration at  $V_{GS}=0.5V$ ,  $V_{DS}=1V$ .

According to Fig.6, with increasing body doping concentration, the electron density of CNTFET increases due to increasing drain current.

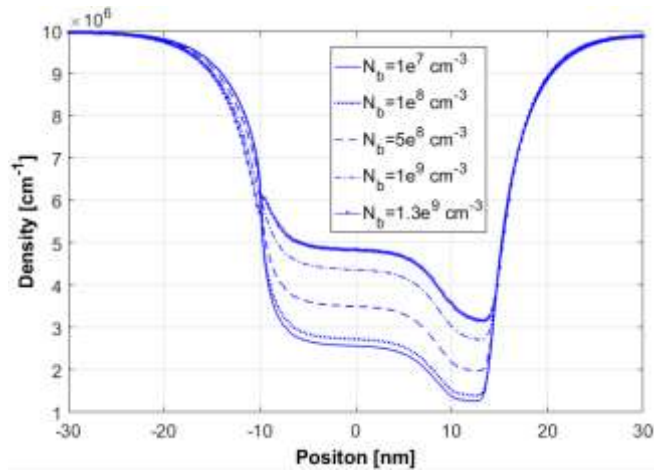


Fig. 6. Electron density of CNTFET for different body doping concentration at  $V_{GS}=0.5V$ ,  $V_{DS}=1V$

In order to investigate body doping concentration effect on the  $I_{ON}/I_{OFF}$  ratio of CNTFET, it is changed in the range of  $N_b=0cm^{-3}$  to  $N_b=1.3 \times e^9 cm^{-3}$  other parameters are constant. The result is shown in Fig.7 with the mentioned parameters with respect to the body doping concentration at  $V_{GS}=0.5V$ ,  $V_{DS}=1V$  for on-state and at  $V_{GS}=0.0V$ ,  $V_{DS}=1V$  for off-state. When the concentration level is below than  $N_b=1.0 \times e^7 cm^{-3}$  Output characteristics change is not noticeable.

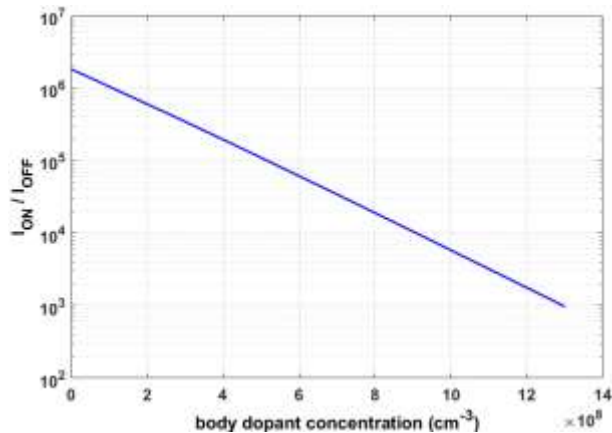


Fig. 7: the  $I_{ON}/I_{OFF}$  ratio of CNTFET for different body doping

The maximum body doping concentration level that we can apply to this CNTFET, is  $N_b=1.3 \times e^9 cm^{-3}$  When the concentration level is higher than this value, simulation results are not agreeable. According to results, with increases in the body doping concentration, on-state and off-state current will increase. But off-state current increment is more than on-state. Therefore the  $I_{ON}/I_{OFF}$  ratio decreases. So, body doping concentration has an inverse relation with the  $I_{ON}/I_{OFF}$  ratio. The height of the barrier is modulated by the body doping concentration. Hence with increasing the body doping level, more carriers flow through the channel and output current will increase.

#### IV. . CONCLUSION

The self-consistent solution of the Schrodinger and Poisson equations has been performed to extract I-V characteristics of carbon nanotube FET with a coaxial gate and zigzag nanotube with (16,0) chirality. For this device, the body doping concentration is increased to  $N_b=1 \times e^9 cm^{-3}$  and results are shown in Fig. 2-7. According to results, with increasing body doping concentration level, the output current, energy profile and electron density of CNTFET will increase and the  $I_{ON}/I_{OFF}$  ratio will decrease.

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